For closed shell of n electrons

$$E_{HF} = 2\sum_{i=1}^{n/2} h_{ii} + \sum_{i=1}^{n/2} J_{ii} + \sum_{i>j}^{n/2-1} \sum_{j=1}^{n/2} (4J_{ij} - 2K_{ij}) + V_{NN}$$

A more compact form is:

$$E_{HF} = 2\sum_{i=1}^{n/2} h_{ii} + \sum_{i=1}^{n/2} \sum_{j=1}^{n/2} (2J_{ij} - K_{ij}) + V_{NN}$$

Which works because $K_{ii} = J_{ii}$ so there is only $1 J_{ii}$ for each i

Applying the Variation Principle

A minimum E is sought by varying the occupied MO's, until E is a minimum. By use of the calculus of variations, and the method of Lagrange undetermined multipliers, the condition of minimum E requires that: $\hat{F}\varphi_i(1) = 2\sum_{i=1}^{n/2} \varepsilon_{ij}\varphi_j(1)$

Where **F** is the Fock Operator, given by: $\hat{F}(1) = \hat{h}(1) + \sum_{j} (2\hat{J}_{j}(1) - \hat{K}_{j}(1))$ Note: **F** is a **one-electron operator!**

How does it get to be 1-electron when H has $1/r_{12}$??

Because, J_j (1) is repulsion an electron with $(\varphi_j)^2$, the electron charge cloud of the jth MO. Not interaction with another electron directly — just the <u>average</u> position.

It is like interacting with a nucleus, except it is a smeared out electron charge

the condition of minimum E requires <u>only</u> that: $\hat{F}\varphi_i(1) = 2\sum_{i=1}^{n/2} \mathcal{E}_{ij}\varphi_j(1)$

Because <u>Any</u> unitary transformation of the set of canonical MOs will yield the same electron density, and therefore the same energy and other properties.

One case of this is most often seen:

$$\hat{F}\varphi_{i}(1) = \mathcal{E}_{i}\varphi_{i}(1)$$

obtained by diagonalizing the matrix:

$$<\!\varphi_{\!i}\,|\,\hat{F}\,|\,\varphi_{j}>$$

These eigenfunctions of the Fock Operator are called the **Canonical Molecular Orbitals**,

The Fock Operator:

The Fock Operator F is an effective Hamiltonian, like the one in Huckel theory, but much more versatile and accurate.

Therefore, the Fock Operator is physically quite transparent and easy to grasp.

$$\hat{F}(1) = \hat{h} + \sum_{j} (2\hat{J}_{j} - \hat{K}_{j})$$

h says that each electron (always called electron 1) has KE and attraction for all the nuclei (the one-electron part, sometime called H^{core}).

 J_j (1) is repulsion of electron 1 with $(\varphi_j)^2$, the electron charge cloud of the jth MO. Not interaction with another electron directly — just the <u>average</u> position

same is true for the Kj

 $-K_j$ reduces this repulsion between the electron in the the jth MO having the same m_s as electron 1.

To see what the J operator (\hat{J}) is, we first write the expression for the J_{ii} integral:

$$J_{ij} = \int \int \varphi_i^*(1) \varphi_j^*(2) \frac{1}{r_{12}} \varphi_i(1) \varphi_j(2) d\tau_1 d\tau_2$$

We may reorder so that all the integration over electron 2 is in the center, revealing the J operator.

$$J_{ij} = \int \varphi_i^*(1) \left[\int \varphi_j^*(2) \frac{1}{r_{12}} \varphi_j(2) d\tau_2 \right] \varphi_i(1) d\tau_1$$

where $\left[\int \varphi_j^*(2) \frac{1}{r_{12}} \varphi_j(2) d\tau_2 \right] = \hat{J}_j(1)$
 $J_{ij} = \int \varphi_i^*(1) \ \hat{J}_j(1) \varphi_i(1) d\tau_1$

Because J_{ij} is just a classical repulsion of two clouds, J_{op} can be thought of as very similar to e^2/r_{12} except that the charge is not a point, it is a cloud shaped like the square of the jth molecular orbital. It is just a function of electron 1 because we integrated over electron 2.

The exchange operator requires special description:

$$K_{ij} = \int \phi_i^*(1) \quad \int \phi_j^*(2)\phi_i(2)\frac{1}{r_{12}}d\tau_2 \quad \phi_j(1)d\tau_1$$

or in Dirac notation:

$$K_{ij} = <\phi_i(1) | \quad <\phi_j(2) | \frac{1}{r_{12}} | \phi_i(2) > | \phi_j(1) >$$

but we want: $K_{ij} = \langle \phi_i (1) | \hat{K}_j | \phi_i (1) \rangle$

 \hat{K}_{j} evidently must turn ϕ_{i} (1) into ϕ_{j} (1) and multiply by

 $<\phi_j(2)|\frac{1}{r_{12}}|\phi_i(2)>$, but it does not know what ϕ_i

it will be operating on;

In all text books I have seen, this operator is just <u>stated</u> to have have the following effect:

$$\hat{K}_{j} | \phi_{i}(1) \rangle = \langle \phi_{j}(2) | \frac{1}{r_{12}} | \phi_{i}(2) \rangle | \phi_{j}(1) \rangle$$

In words: $turn |\phi_i| (1) > into |\phi_j| (1) >$

and multiply by the number: $\langle \phi_j(2) | \frac{1}{r_{12}} | \phi_i(2) \rangle$ This is just the overlap integral, which would be zero, except it is weighted by $1/r_{12}$ in the integral. An explicit operator that will accomplish this is :

$$\hat{K}_{j} = \langle \varphi_{j}(2) | \frac{1}{r_{12}} \left\{ \sum_{\mathbf{m}} | \varphi_{\mathbf{m}}(2) > | \varphi_{j}(1) > \langle \varphi_{\mathbf{m}}(1) | \right\}$$

then

No matter what value of i, K_j gives 0, unless m=i; then we get:

$$\hat{K}_{j} | \varphi_{i}(1) \rangle = \langle \varphi_{j}(2) | \frac{1}{r_{12}} \varphi_{i}(2) \rangle | \varphi_{j}(1) \rangle$$



a Raising Operator is just a sum of transition operators

Transition Density == $\phi_i(1)\phi_j^*(1)$ i.e. a product of two different functions

Transition Density operator = $|\phi_i| > \langle \phi_j| |$

changes $|\phi_j(1)\rangle$ into $|\phi_i(1)\rangle$

Raising operator

 $a = \sum_{\mathbf{m}} |\phi_{\mathbf{m}+1}(1) > \langle \phi_{\mathbf{m}}(1)|$ $a |\phi_{n}(1) > = |\phi_{n+1}(1) >$ $\sum_{\mathbf{m}} |\phi_{\mathbf{m}+1}(1) > \langle \phi_{\mathbf{m}}(1)| |\phi_{n}(1) >$ $= 0 \text{ unless } \mathbf{m} = \mathbf{n}, \text{ which gives } |\phi_{n+1}(1) >$

Lowering operator

$$a^{\dagger} = \sum_{\mathbf{m}} |\phi_{\mathbf{m}-1}(1) > \langle \phi_{\mathbf{m}}(1)|$$

$$a^{\dagger} |\phi_{n}(1) > = |\phi_{n-1}(1) >$$

$$\sum_{\mathbf{m}} |\phi_{\mathbf{m}-1}(1) > \langle \phi_{\mathbf{m}}(1)| |\phi_{n}(1) >$$

$$= 0 \text{ unless } \mathbf{m} = \mathbf{n}, \text{ which gives } |\phi_{n-1}(1) >$$

Raising operator operating on $|\phi_n(1)\rangle$ gives 0 unless m=n

$$a = \sum_{\mathbf{m}} |\phi_{\mathbf{m}+1}(1) > \langle \phi_{\mathbf{m}}(1)|$$

$$a | \phi_n(1) \rangle = | \phi_{n+1}(1) \rangle$$

$$\sum_{\mathbf{m}} | \phi_{\mathbf{m}+1}(1) \rangle \langle \phi_{\mathbf{m}}(1) | | \phi_n(1) \rangle$$

= 0 unless m = n, which gives $|\phi_{n+1}(1)\rangle$

Lowering operator

$$a^{\dagger} = \sum_{\mathbf{m}} |\phi_{\mathbf{m}-1}(1) > \langle \phi_{\mathbf{m}}(1)|$$

$$a^{\dagger} |\phi_{n}(1) > = |\phi_{n-1}(1) >$$

$$\sum_{\mathbf{m}} |\phi_{\mathbf{m}-1}(1) > \langle \phi_{\mathbf{m}}(1)| |\phi_{n}(1) >$$

= 0 unless m = n, which gives $|\phi_{n-1}(1) >$

What does $|i\rangle\langle j| + |j\rangle\langle k| + |k\rangle\langle i|$ do ? i,j,k are Cartesian unit vectors in x,y,z space

Orbital Energy

If $F\varphi_i = \varepsilon_i \varphi_i$ for each MO that is occupied, the energy of the Slater determinant is **minimized**, and

 \mathcal{E}_{i} is called the "ith MO energy" = $\langle \varphi_{i} | F | \varphi_{i} \rangle$. Where, $\langle \varphi_{i} | F | \varphi_{i} \rangle = h_{ii} + \sum_{j} (2Jij - Ki_{j})$

 $\varepsilon_i = h_{ii} + J_{ii} + sum of 2J_{ij} - K_{ij}$ for all j MOs for j not equal to i

Physical Interpretation of "MO Energy"

Pretend you are an electron: what you have and see is your energy. (you are in MO i)

E_i is your:
Kinetic energy

+ your attraction coulombic attraction to all the nuclei,

- + the repulsion energy of your cloud of negative charge with each of the other clouds j
- the repulsion energy of the "charges" given by the product of your orbital (i) times those orbitals j with the same m_s that you have.

 $\varepsilon_i = h_{ii} + J_{ii} + sum of 2J_{ij} - K_{ij}$ for all j MOs for j not equal to i

E_i = h_{ii} + J_{ii} + sum of 2J_{ij} –K_{ij} for all j MOs for j not equal to i Is the Sum of MO Energies the Total HF Energy? The Sum of MO Energies is <u>NOT</u> the Total HF Energy

The sum of MO energies is

$$\sum_{i=1}^{n/2} 2\varepsilon_i = \sum_{i=1}^{n/2} 2h_{ii} + \sum_{i=1}^{n/2} \sum_{j=1}^{n/2} (4J_{ij} - 2K_{ij})$$

but,
$$E_{HF} = \sum_{i=1}^{n/2} 2h_{ii} + \sum_{i=1}^{n/2} \sum_{j=1}^{n/2} (2J_{ij} - K_{ij})$$

The sum of MO energies counts the electron repulsion TWICE.

$$\sum_{i=1}^{n/2} 2\varepsilon_i = \sum_{i=1}^{n/2} 2h_{ii} + \sum_{i=1}^{n/2} \sum_{j=1}^{n/2} (4J_{ij} - 2K_{ij})$$

So, $E_{HF} = \sum_{i=1}^{n/2} 2\varepsilon_i - \sum_{i=1}^{n/2} \sum_{j=1}^{n/2} (2J_{ij} - K_{ij})$

Which is the sum of MO energies minus the total electron repulsion



Koopman's Theorem:

The ith "MO Energy" = minus the ionization potential for removing an electron from the ith MO

Remove one of the electrons in MO A

$$\underline{E}_{cation} = \mathbf{h}_{AA} + 2\mathbf{h}_{BB} + \mathbf{J}_{BB} + \mathbf{2J}_{BA} - \mathbf{K}_{BA}$$
$$\mathbf{E}_{difference} = \mathbf{h}_{AA} + \mathbf{J}_{AA} + 2\mathbf{J}_{BA} - \mathbf{K}_{BA} = \mathbf{E}_{A}$$
$$= \mathbf{The ionization energy}$$

If no relaxation due to changed electron density



Virtual (i.e., unoccupied orbitals are different!

Add an electron to MO C

The ith "MO Energy" is for an electron that "sees" <u>ALL of the</u> <u>occupied MO electrons</u>

$$E_{neutral} = 2h_{AA} + 2h_{BB} + J_{AA} + J_{BB} + 4J_{BA} - 2K_{BA}$$

$$E_{anion} = 2h_{AA} + 2h_{BB} + J_{AA} + J_{BB} + 4J_{BA} - 2K_{BA} + h_{CC} + 2J_{CA} - K_{CA} + 2J_{CB} - K_{CB}$$

$$E_{difference} = h_{CC} + 2J_{CA} - K_{CA} + 2J_{CB} - K_{CB} = \mathcal{E}_{C}$$

$$= minus the Electron Affinity$$
If no relaxation due to changed electron density